

## Probing DE w/ Cluster Counts

Recall FRW metric:

$$ds^2 = dt^2 - a^2(t) \left[ \frac{dr^2}{(1-kr^2)} + r^2(\delta\theta^2 + \sin^2\theta d\phi^2) \right]$$

## Volume Element

- If we have a set of objects (e.g., galaxies, clusters, dark matter halos) whose spatial number density we know / can predict, then counting them provides another cosmological test. We'll apply this esp. to galaxy cluster counts.

~~Proper area~~ A proper area  $dA$  at redshift  $z$  subtends solid angle  $d\Omega$  at the origin given by  $dA = a(t_0) r d\theta d\phi \sin\theta d\Omega$

$$= a_0^2 r^2 d\Omega = \frac{a_0^2 r^2 d\Omega}{(1+z)^2}$$

- The rate of proper displacement w/z along a light ray is

$$dl = c dt = \frac{da}{\dot{a}} = \frac{dz}{(1+z)} \frac{a}{\dot{a}} = \frac{dz}{H(z)(1+z)}$$

= linear depth of a sample of objects in the redshift interval  $(z, z+\delta z)$

- Proper volume element of the sample is then

$$dV_{\text{prop}} = dA dl = \frac{a_0^2 r^2(z)}{H_0 E(z) (1+z)^3} dl dz \quad \text{where } H(z) = H_0 E(z)$$

- If  $n_{\text{prop}}(z)$  is the proper number density of objects, then the counts per unit redshift and solid angle are

$$\frac{d^2 N}{dz d\Omega} = n_{\text{prop}}(z) \frac{d^2 V_{\text{prop}}}{dz d\Omega} = n_{\text{prop}}(z) \frac{a_0^2 r^4(z)}{H_0 E(z) (1+z)^3}$$

- If objects are conserved (neither created nor destroyed), then

$$n_{\text{prop}}(z) = n_0 (1+z)^3 \Rightarrow \frac{d^2N}{dz d\Omega} = \frac{n_0 r^2(z)}{H(z)}$$

- Convenient to define the comoving # density,

$$n_c(z) = \frac{n_{\text{prop}}(z)}{(1+z)^3} = \text{constant if objects are conserved}$$

= # density per comoving volume,

where the comoving volume element is defined by

$$\frac{d^2V_c}{dz d\Omega} = \frac{r^2(z)}{H(z)} = \frac{r^2(z)}{H_0 E(z)} = \left( \frac{d^2V_{\text{prop}}}{dz d\Omega} \right) (1+z)^3$$

so that the counts are given by

$$\frac{d^2N}{dz d\Omega} = n_c(z) \frac{d^2V_c}{dz d\Omega} = n_c(z) \frac{r^2(z)}{H(z)}$$

- See Fig. 2 right-hand panel of Rio lectures for example & Fig. 9 Rio

- We will apply this to clusters, for which theory predicts  $n(z)$ .
- See e.g. Vikhlinin et al results

$$\frac{\partial^2 V}{\partial z \partial \bar{z}} \text{ scaling: } \sim \frac{r^2}{H(z)}$$

~~$$x = \int \frac{dz}{H(z)}$$~~

$$r = S_e(x) \approx x \text{ at low } z \ll 1$$

$$H(z) = H_0 \left[ \Omega_m (1+z)^3 + \Omega_{DE} (1+z)^{3(1+w)} \right]^{1/2}$$

Expand:  
for  $z \ll 1$

$$H(z) \approx H_0 \left( 1 + (1+q_0)z \right)$$

$$q_0 = \frac{\Omega_m}{2} + \frac{(1-\Omega_m)}{2} (1+3w) \quad \text{for flat}$$

$$X \approx \frac{1}{H_0} \left[ z - (1+q_0) \frac{z^2}{2} \right]$$

$$\text{To lowest order, } \frac{\partial^2 V}{\partial z \partial \bar{z}} \sim \frac{r^2(z)}{H(z)} \sim \frac{X'(z)}{H(z)}$$

$$\frac{\partial^2 V}{\partial z \partial \bar{z}} \propto \frac{\left[ z - (1+q_0) \frac{z^2}{2} \right]^2}{1 + (1+q_0)z} \sim z^{-2} \frac{\left[ 1 - (1+q_0) \frac{z}{2} \right]^2}{1 + (1+q_0)z}$$

$$\frac{\partial^2 V}{\partial z \partial \bar{z}} \underset{z \ll 1}{\propto} z^2 \left( 1 + (1+\epsilon_0) \frac{z^2}{4} - (1+\epsilon_0) z \right) \left( 1 - (1+\epsilon_0) z \right)$$

$$\propto z^2 \left[ 1 - (1+\epsilon_0)z - (1+\epsilon_0)z + O(z^2) \right]$$

$$\propto z^2 \left[ 1 - 2(1+\epsilon_0)z \right]$$

For  $\omega_m = 0.2$ ,  $\epsilon_0 = 0.1 + 0.4(1+3\omega)$

$$\omega = -1 : \epsilon_0 = 0.1 + 0.4(1-3)$$

$$= 0.1 - 0.8 = -0.7$$

$$(1+\epsilon_0 = 0.3)$$

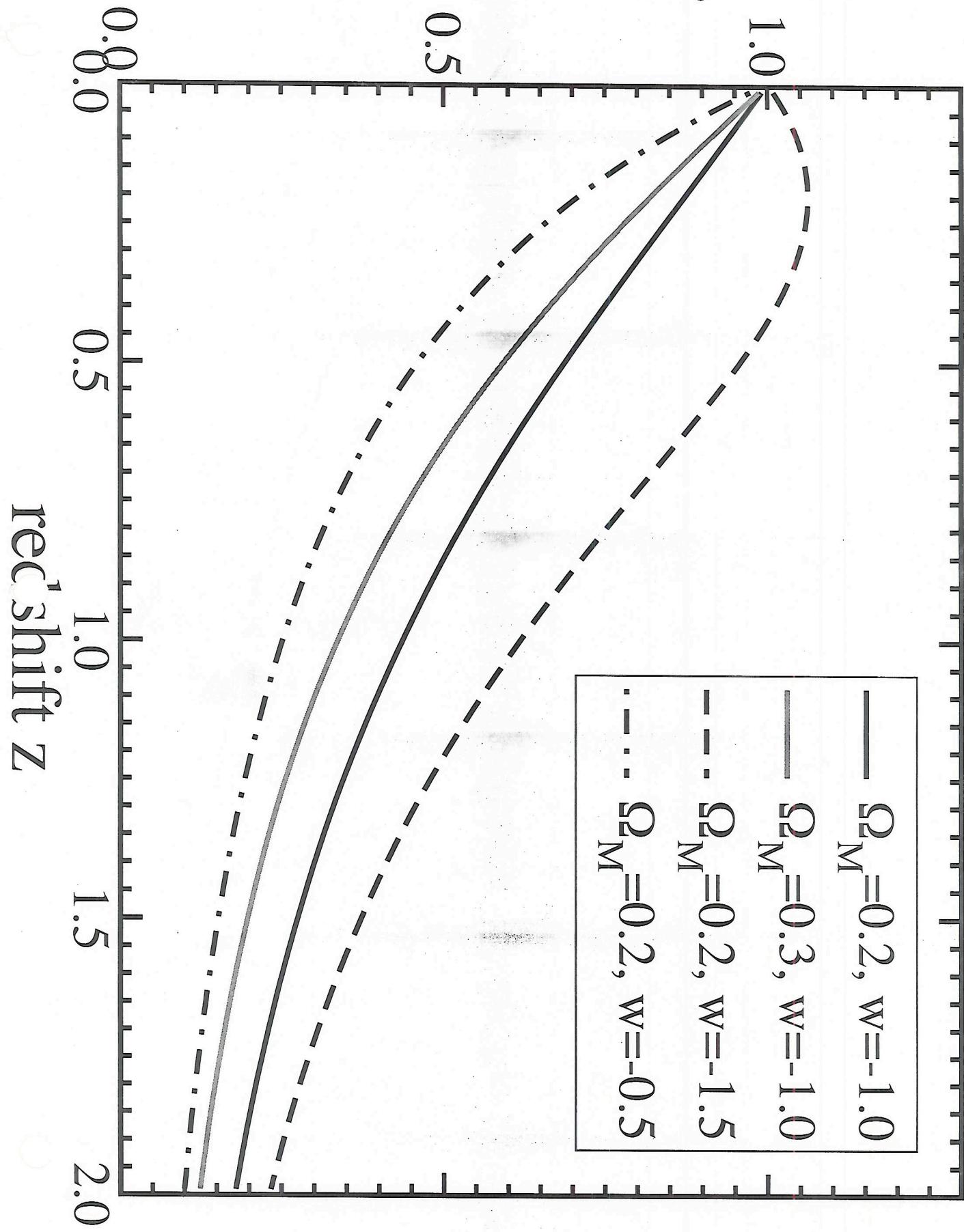
$$\omega = -0.5 : \epsilon_0 = 0.1 + 0.4(1-1.5)$$

$$= 0.1 - 0.2 = -0.1$$

$$(1+\epsilon_0 = 0.9)$$

$$V_{01} \propto \begin{cases} z^2 \left[ 1 - 0.6z \right], & \omega = -1 \\ z^2 \left[ 1 - 1.8z \right], & \omega = -0.5 \end{cases}$$

$$(d^2V/d\Omega dz) \times (H_0^3 / z^2)$$



## Counting Clusters

$$\frac{d^2 N(z)}{dz d\Omega} = \frac{r^2(z)}{H(z)} \int_0^\infty f(O, z) d\Omega \int_0^\infty p(O | M, z) \frac{dn(z)}{dM} dM$$

selection fr. for  
 observable  $O$ , e.g.  
 optical richness, & say luminosity,  
~~flux~~ flux, bc mass  
 not constant  
 with redshift

mass-observable  
 relation  
 theoretical  
 mass  
 fn.

Q: Vikhlinin et al.

recall

~~$D(z) = r(z)(1+z)^{-1}$~~

$D(z) = r(z)(1+z)$

Hal. Mass  
 function,  
 predicted by  
 N-body  
 simulations

appro. by  
 Press-Schechter Fn.

$\mathbf{g}$   
 encodes DE-dependence  
 of perturbation  
 growth